Continuous-Time Fourier Transform

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From Periodic to Aperiodic

Previously we have focused on Fourier representations of periodic signals: e.g., sounds, waves, music, ...

However, most real-world signals are not periodic. None are truly periodic since they do not have infinite duration!

Today: generalizing Fourier representations to include aperiodic signals.
How can we represent an aperiodic signal as a sum of sinusoids?
Fourier Representations of Aperiodic Signals

How can we represent an aperiodic signal as a sum of sinusoids?

Strategy: make a periodic version of \( f(t) \) by summing shifted copies:

\[
f_p(t) = \sum_{m=-\infty}^{\infty} f(t - mT)
\]

Since \( f_p(t) \) is periodic, it has a Fourier series (which depends on \( T \)).

Find Fourier series coefficients \( F_p[k] \) and take the limit of \( F_p[k] \) as \( T \to \infty \).

As \( T \to \infty \), \( f_p(t) \to f(t) \) and Fourier series will approach Fourier transform.
Fourier Representations of Aperiodic Signals

Example.

Strategy: make a periodic version of $f(t)$ by summing shifted copies:

$$f_p(t) = \sum_{m=-\infty}^{\infty} f(t - mT)$$

Calculate the Fourier series coefficients $F_p[k]$:

$$F_p[k] = \frac{1}{T} \int_{-S}^{S} e^{-j \frac{2\pi k}{T} t} \, dt = \frac{1}{T} \left. \frac{e^{-j \frac{2\pi k}{T} t}}{-j \frac{2\pi k}{T}} \right|_{-S}^{S} = \frac{2 \sin \left( \frac{2\pi k}{T} S \right)}{T \left( \frac{2\pi k}{T} \right)}$$
Calculate the Fourier series coefficients $F_p[k]$:

$$F_p[k] = \frac{1}{T} \int_{-S}^{S} e^{-j \frac{2\pi}{T} kt} dt = \frac{1}{T} \left[ \frac{e^{-j \frac{2\pi}{T} kt}}{-j \frac{2\pi k}{T}} \right]_{-S}^{S} = \frac{2 \sin \left( \frac{2\pi k}{T} S \right)}{T \left( \frac{2\pi k}{T} \right)}$$

Plot the resulting Fourier coefficients when $S=1$ and $T=8$.

What happens if you double the period $T$?
Calculate the Fourier series coefficients \( F_p[k] \):

\[
F_p[k] = \frac{1}{T} \int_{-S}^{S} e^{-j \frac{2\pi}{T} kt} \, dt = \frac{1}{T} e^{-j \frac{2\pi}{T} k t} \bigg|_{-S}^{S} = \frac{2 \sin \left( \frac{2\pi k}{T} S \right)}{T \left( \frac{2\pi k}{T} \right)}
\]

Plot the resulting Fourier coefficients when \( S=1 \) and \( T=8 \).

What happens if you double the period \( T \)? Plot with \( S=1 \) and \( T=16 \).

There are twice as many samples per period of the sin function. (The red samples are at new intermediate frequencies.) The amplitude is halved.
Fourier Representations of Aperiodic Signals

Define a new function $F(\omega)$ where $\omega = k\omega_o = \frac{2\pi k}{T}$.

$$TF_p[k] = 2\sin\left(\frac{2\pi k}{T}S\right) = 2\frac{\sin(\omega S)}{\omega}\bigg|_{\omega = \frac{2\pi k}{T}} = F(\omega)|_{\omega = \frac{2\pi k}{T}}$$

Then $TF_p[k]$ represents samples of $F(\omega)$ with increasing resolution in $\omega$.

$S=1$ and $T=8$:

$S=1$ and $T=16$:

$S=1$ and $T=32$:

The discrete function $TF_p[k]$ is a sampled version of the function $F(\omega)$. 
Fourier Representations of Aperiodic Signals

How do we find $F(\omega)$?

We started with a periodic version of $f(t)$ that has a FS:

$$f_p(t) \overset{{\text{CTFS}}}{\rightarrow} F_p[k]$$

Take the limit as $T \rightarrow \infty$ of both sides:

$$\lim_{T \rightarrow \infty} f_p(t) = f(t) \overset{{\text{CTFS}}}{\rightarrow} \lim_{T \rightarrow \infty} F_p[k] = \lim_{T \rightarrow \infty} \left[ \frac{1}{T} F(\omega) \right]_{\omega = \frac{2\pi k}{T}}$$

Try

$$F(\omega) = \lim_{T \rightarrow \infty} T F_p[k] \bigg|_{k = \frac{T}{2\pi} \omega}$$

$$= \lim_{T \rightarrow \infty} T \left[ \frac{1}{T} \int_T f_p(t) e^{-j \frac{2\pi k}{T} t} dt \right] \bigg|_{k = \frac{T}{2\pi} \omega}$$

$$= \lim_{T \rightarrow \infty} \int_T f_p(t) e^{-j \omega t} dt$$

$$F(\omega) = \int_{-\infty}^{\infty} f(t) e^{-j \omega t} dt$$
Fourier Representations of Aperiodic Signals

We can reconstruct $f(t)$ from $F(\omega)$ using a Riemann sum.

$$f(t) = \lim_{T \to \infty} f_p(t) = \lim_{T \to \infty} \sum_k F_p[k] e^{j \frac{2\pi}{T} kt}$$

$$= \lim_{T \to \infty} \left( \frac{1}{2\pi} \right) \sum_k T F_p[k] e^{j \frac{2\pi}{T} kt} \left( \frac{2\pi}{T} \right) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega) e^{j \omega t} d\omega$$

$TF_p[k] = F(\omega)$

$S=1$ and $T=8$: $\omega = \frac{2\pi k}{T}$

$S=1$ and $T=16$: $\omega = \frac{2\pi k}{T}$

$S=1$ and $T=32$: $\omega = \frac{2\pi k}{T}$
Continuous-Time Fourier Representations

Fourier series and transforms are similar: both represent signals by their frequency content.

Continuous-Time Fourier Series

\[ F[k] = \frac{1}{T} \int_{T} f(t) e^{-jk\omega_0 t} dt \quad \text{analysis equation} \]

\[ f(t) = f(t + T) = \sum_{k=-\infty}^{\infty} F[k] e^{jk\omega_0 t} \quad \text{synthesis equation} \]

where \( \omega_0 = \frac{2\pi}{T} \)

Continuous-Time Fourier Transform

\[ F(\omega) = \int_{-\infty}^{\infty} f(t) e^{-j\omega t} dt \quad \text{analysis equation} \]

\[ f(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega) e^{j\omega t} d\omega \quad \text{synthesis equation} \]
Continuous-Time Fourier Representations

All of the information in a periodic signal is contained in one period. The information in an aperiodic signal is spread across all time.

Continuous-Time Fourier Series

$$F[k] = \frac{1}{T} \int_T f(t) e^{-jk\omega_0 t} dt$$  \hspace{1cm} \text{analysis equation}

$$f(t) = f(t + T) = \sum_{k=-\infty}^{\infty} F[k] e^{jk\omega_0 t}$$ \hspace{1cm} \text{synthesis equation}

where \( \omega_0 = \frac{2\pi}{T} \)

Continuous-Time Fourier Transform

$$F(\omega) = \int_{-\infty}^{\infty} f(t) e^{-j\omega t} dt$$ \hspace{1cm} \text{analysis equation}

$$f(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega) e^{j\omega t} d\omega$$ \hspace{1cm} \text{synthesis equation}
Continuous-Time Fourier Representations

Periodic signals can be synthesized from a discrete set of harmonics. Aperiodic signals generally require all possible frequencies.

Continuous-Time Fourier Series

\[ F[k] = \frac{1}{T} \int_T f(t) e^{-jk\omega_0 t} dt \]  
\[ f(t) = f(t + T) = \sum_{k=-\infty}^{\infty} F[k] e^{jk\omega_0 t} \]

where \( \omega_0 = \frac{2\pi}{T} \)

Continuous-Time Fourier Transform

\[ F(\omega) = \int_{-\infty}^{\infty} f(t) e^{-j\omega t} dt \]
\[ f(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega) e^{j\omega t} d\omega \]
Continuous-Time Fourier Representations

Harmonic frequencies $k\omega_o$ are samples of continuous frequency $\omega$.

Continuous-Time Fourier Series

$$F[k] = \frac{1}{T} \int_T f(t)e^{-jk\omega_o t} dt$$  \hspace{1cm} \text{analysis equation}

$$f(t) = f(t + T) = \sum_{k=-\infty}^{\infty} F[k]e^{jk\omega_o t} \hspace{1cm} \text{synthesis equation}$$  \hspace{1cm} \text{where } \omega_o = \frac{2\pi}{T}

Continuous-Time Fourier Transform

$$F(\omega) = \int_{-\infty}^{\infty} f(t)e^{-j\omega t} dt$$  \hspace{1cm} \text{analysis equation}

$$f(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega)e^{j\omega t} d\omega$$ \hspace{1cm} \text{synthesis equation}
Example

Find the Fourier Transform (FT) of a rectangular pulse:

\[
f(t) = \begin{cases} 
1 & -1 < t < 1 \\
0 & \text{otherwise}
\end{cases}
\]

\[
F(\omega) = \int_{-\infty}^{\infty} f(t)e^{-j\omega t} dt = \int_{-1}^{1} e^{-j\omega t} dt = \frac{e^{-j\omega t}}{-j\omega} \bigg|_{-1}^{1} = 2 \frac{\sin \omega}{\omega}
\]

\[
F(\omega)
\]

\[
-4\pi \quad -2\pi \quad 0 \quad 2\pi \quad 4\pi
\]

\[
f(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega)e^{j\omega t} d\omega
\]

A square pulse contains (almost) all frequencies \(\omega\) (missing just \(\pi, 2\pi, \ldots\)).
Properties of Fourier Transforms

Fourier transforms offer an alternative view of an aperiodic signal.

A signal and its Fourier transform contain exactly the same information, but some information is more easily seen in one domain than in the other.

There are many properties of Fourier transforms. These properties summarize systematic relations between time and frequency representations.
Properties of Fourier Transforms

Time delay maps to linear phase delay of the Fourier transform.

If

\[ f(t) \xrightarrow{\text{FT}} F(\omega) \]

then

\[ f(t - \tau) \xrightarrow{\text{FT}} e^{-j\omega \tau} F(\omega) \]

\[
F(\omega) = \int_{-\infty}^{\infty} f(t)e^{-j\omega t} dt
\]

\[
G(\omega) = \int_{-\infty}^{\infty} f(t - \tau)e^{-j\omega t} dt
\]

Let \( u = t - \tau \) (and therefore \( du = dt \) since \( \tau \) is a constant)

\[
G(\omega) = \int_{-\infty}^{\infty} f(u)e^{-j\omega(u+\tau)} du = e^{-j\omega \tau} \int_{-\infty}^{\infty} f(u)e^{-j\omega u} du = e^{-j\omega \tau} F(\omega)
\]

The angle of \( e^{-j\omega \tau} = -\omega \tau \).

Why does time delay change phase by an amount proportional to frequency?
Properties of Fourier Transforms

Why does time delay change phase by an amount proportional to frequency?

Doubling the frequency of a sinusoid doubles the change in phase associated with a given time delay.
Properties of Fourier Transforms

Scaling time.

Consider the following signal and its Fourier transform.

Time representation:

\[ f_1(t) = \begin{cases} 1 & \text{for } -1 < t < 1 \\ 0 & \text{otherwise} \end{cases} \]

Frequency representation:

\[ F_1(\omega) = \frac{2 \sin \omega}{\omega} \]

How would these functions scale if time were stretched?
Check Yourself

Signal $f_2(t)$ and its Fourier transform $F_2(\omega)$ are shown below.

Which of the following is true?

1. $b = 2$ and $\omega_0 = \pi/2$
2. $b = 2$ and $\omega_0 = 2\pi$
3. $b = 4$ and $\omega_0 = \pi/2$
4. $b = 4$ and $\omega_0 = 2\pi$
5. none of the above
Check Yourself

Find the Fourier transform.

\[ F_2(\omega) = \int_{-2}^{2} e^{-j\omega t} \, dt = \frac{e^{-j\omega t}}{j\omega} \bigg|_{-2}^{2} = \frac{2 \sin 2\omega}{\omega} = \frac{4 \sin 2\omega}{2\omega} \]

\[ F_1(\omega) = \frac{2 \sin \omega}{\omega} \]

\[ F_2(\omega) = \frac{4 \sin 2\omega}{2\omega} \]

Stretching time compresses frequency.
Signal $f_2(t)$ and its Fourier transform $F_2(\omega)$ are shown below.

Which of the following is true? 3

1. $b = 2$ and $\omega_0 = \pi/2$
2. $b = 2$ and $\omega_0 = 2\pi$
3. $b = 4$ and $\omega_0 = \pi/2$
4. $b = 4$ and $\omega_0 = 2\pi$
5. none of the above
Properties of Fourier Transforms

Find a general scaling rule.

Let \( f_2(t) = f_1(at) \) where \( a > 0 \).

\[
F_2(\omega) = \int_{-\infty}^{\infty} f_1(at) e^{-j\omega t} dt
\]

Let \( \tau = at \). Then \( d\tau = a \, dt \).

\[
F_2(\omega) = \int_{-\infty}^{\infty} f_1(\tau) e^{-j\omega \tau/a} \frac{1}{a} d\tau = \frac{1}{a} F_1 \left( \frac{\omega}{a} \right)
\]

Stretching time compresses frequency and increases amplitude (preserving area).
The value of $F(\omega)$ at $\omega = 0$ is the integral of $f(t)$ over time $t$.

$$F(\omega)\big|_{\omega=0} = \int_{-\infty}^{\infty} f(t)e^{-j\omega t} dt = \int_{-\infty}^{\infty} f(t)e^{-j0t} dt = \int_{-\infty}^{\infty} f(t) dt$$

$F_1(\omega) = \frac{2\sin \omega}{\omega}$
Moments

The value of $f(0)$ is the integral of $F(\omega)$ divided by $2\pi$.

$$f(0) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega) e^{j\omega t} d\omega = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega) d\omega$$

![Graph showing $f_1(t)$ and $F_1(\omega)$]
Moments

The value of $f(0)$ is the integral of $F(\omega)$ divided by $2\pi$.

$$f(0) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega) e^{j\omega t} d\omega = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega) d\omega$$

![Graph of $f_1(t)$ and $F_1(\omega)$](image)
Stretching Time

Stretching time compresses frequency and increases amplitude (preserving area).

\[ f_1(t) = \begin{cases} 1 & \text{for } -1 < t < 1 \\ 0 & \text{otherwise} \end{cases} \]

\[ F_1(\omega) = \frac{2 \sin \omega}{\omega} \]

![Graph showing the function and its Fourier transform](image)
Compressing Time to the Limit

Alternatively, we could compress time while keeping area = 1.

\[ f(t) \]

\[ F(\omega) = \frac{\sin \omega/2}{\omega/2} \]

In the limit, the pulse has zero width but area 1!

We represent this limit with the delta (or impulse) function: \( \delta(t) \).
Although physically unrealizable, the impulse (a.k.a. Dirac delta) function is useful as a mathematically tractable approximation to a very brief signal.

Example 1: Find the Fourier transform of a unit impulse function.

\[
F(\omega) = \int_{-\infty}^{\infty} \delta(t)e^{-j\omega t} dt
\]

Since \( \delta(t) \) is zero except near \( t=0 \), only values of \( e^{-j\omega t} \) near \( t=0 \) are important. Because \( e^{-j\omega t} \) is a smooth function of \( t \), \( e^{-j\omega t} \) can be replaced by \( e^{-j\omega 0} \):

\[
F(\omega) = \int_{-\infty}^{\infty} \delta(t)e^{-j\omega 0} dt = 1
\]

This matches our previous result which was based explicitly on a limit. Here the limit is implicit.
Math With Impulses

Although physically unrealizable, the impulse (a.k.a. Dirac delta) function is useful as a mathematically tractable approximation to a very brief signal.

Example 2: Find the function whose Fourier transform is an impulse.

\[
f(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \delta(\omega)e^{j\omega t} d\omega = \frac{1}{2\pi} \int_{-\infty}^{\infty} \delta(\omega)e^{j0t} d\omega = \frac{1}{2\pi}
\]

\[
1 \overset{\text{CTFT}}{\Rightarrow} 2\pi \delta(\omega)
\]

Notice the similarity to the previous result:

\[
\delta(t) \overset{\text{CTFT}}{\Rightarrow} 1
\]

These relations are duals of each other.

- A constant in time consists of a single frequency at \( \omega = 0 \).
- An impulse in time contains components at all frequencies.
Math With Impulses

Although physically unrealizable, the impulse (a.k.a. Dirac delta) function is useful as a mathematically tractable approximation to a very brief signal.

Example 3: Find the function whose Fourier transform is a shifted impulse.

\[
f(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \delta(\omega-\omega_o)e^{j\omega t} d\omega = \frac{1}{2\pi} \int_{-\infty}^{\infty} \delta(\omega')e^{j(\omega'+\omega_o)t} d\omega'
\]

\[
= \frac{1}{2\pi} e^{j\omega_o t} \int_{-\infty}^{\infty} \delta(\omega') e^{j\omega' t} d\omega' = \frac{1}{2\pi} e^{j\omega_o t}
\]

\[
e^{j\omega_o t} \overset{\text{CTFT}}{\Rightarrow} 2\pi \delta(\omega-\omega_o)
\]

Use this result to relate Fourier series to Fourier transforms.
Relation Between Fourier Series and Fourier Transforms

If a periodic signal $f(t) = f(t + T)$ has a Fourier series representation, then it can also be represented by an equivalent Fourier transform.

$$e^{j\omega_0 t} \overset{\text{FT}}{\longrightarrow} 2\pi \delta(\omega - \omega_0)$$

$$f(t) = f(t + T) = \sum_{k=-\infty}^{\infty} F[k] e^{j\frac{2\pi}{T}kt} \quad \text{CTFS} \quad \longleftrightarrow \quad F[k]$$

$$f(t) = f(t + T) = \sum_{k=-\infty}^{\infty} F[k] e^{j\frac{2\pi}{T}kt} \quad \text{CTFT} \quad \longleftrightarrow \quad \sum_{k=-\infty}^{\infty} 2\pi F[k] \delta\left(\omega - \frac{2\pi}{T}k\right)$$

Each term in the Fourier series is replaced by an impulse in the Fourier transform.
Relation between Fourier Transform and Fourier Series

Each Fourier series term is replaced by an impulse in the Fourier transform.

\[
f(t) = \sum_{k=-\infty}^{\infty} f_p(t - kT)
\]

\[
F[k]
\]

\[
F(\omega) = \sum_{k=-\infty}^{\infty} 2\pi F[k] \delta(\omega - k\frac{2\pi}{T})
\]
Summary

Fourier series and transforms are similar:
both represent signals by their frequency content.

Continuous-Time Fourier Series

\[ F[k] = \frac{1}{T} \int_{T} f(t)e^{-jk\omega_0 t} dt \]  
\[ f(t) = f(t + T) = \sum_{k=-\infty}^{\infty} F[k]e^{jk\omega_0 t} \]

Continuous-Time Fourier Transform

\[ F(\omega) = \int_{-\infty}^{\infty} f(t)e^{-j\omega t} dt \]  
\[ f(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega)e^{j\omega t} d\omega \]

Next time: Fourier Transform for discrete-time signals.