

6.003: Signal Processing

Fourier Series – Complex Exponential Form

- complex numbers
- complex exponentials and their relation to sinusoids
- complex exponential form of Fourier series
- delay property of Fourier series

September 16, 2021

Fourier Series

We have previously represented signals as weighted sums of sinusoids.

Synthesis Equation

$$f(t) = c_0 + \sum_{k=1}^{\infty} c_k \cos(k\omega_o t) + \sum_{k=1}^{\infty} d_k \sin(k\omega_o t) \quad \text{where } \omega_o = \frac{2\pi}{T}$$

Analysis Equations

$$c_0 = \frac{1}{T} \int_T f(t) dt$$

$$c_k = \frac{2}{T} \int_T f(t) \cos(k\omega_o t) dt$$

$$d_k = \frac{2}{T} \int_T f(t) \sin(k\omega_o t) dt$$

Today: Simplifying the math with complex numbers.

Simplifying Math By Using Complex Numbers

Complex numbers simplify thinking about roots of numbers / polynomials:

- all numbers have two square roots, three cube roots, etc.
- all polynomials of order n have n roots (some of which may be repeated).

→ much simpler than the rules that govern purely real-valued formulations. For example, a cubic equation with real-valued coefficients might have 1 or 3 real-valued roots; a quartic equation might have 0, 2, or 4.

Complex exponentials simplify thinking about trigonometric functions (Euler's formula, Leonhard Euler, 1748):

$$e^{j\theta} = \cos \theta + j \sin \theta$$

where $j = \sqrt{-1}$.

This single equation virtually eliminates our need for trig tables. Richard Feynman called this "the most remarkable formula in mathematics."

Note that we will normally use j (instead of i) to represent $\sqrt{-1}$.

Where Does Euler's Formula Come From?

Euler showed the relation between complex exponentials and sinusoids by solving the following differential equation two ways.

$$\frac{d^2 f(\theta)}{d\theta^2} + f(\theta) = 0$$

$$\text{let } f_1(\theta) = A \cos(\alpha\theta) + B \sin(\beta\theta)$$

$$\frac{df_1(\theta)}{d\theta} = -\alpha A \sin(\alpha\theta) + \beta B \cos(\beta\theta)$$

$$\frac{d^2 f_1(\theta)}{d\theta^2} = -\alpha^2 A \cos(\alpha\theta) - \beta^2 B \sin(\beta\theta)$$

$$\alpha^2 = \beta^2 = 1$$

$$f_1(\theta) = A \cos \theta + B' \sin \theta$$

$$\text{Let } f_2(\theta) = C e^{\gamma\theta}$$

$$\frac{df_2(\theta)}{d\theta} = \gamma C e^{\gamma\theta}$$

$$\frac{d^2 f_2(\theta)}{d\theta^2} = \gamma^2 C e^{\gamma\theta}$$

$$\gamma^2 = -1$$

$$f_2(\theta) = C e^{\pm j\theta}$$

If we arbitrarily take $f_2(\theta) = e^{j\theta}$, $f_2(0) = 1$ and $f_2'(0) = j$.

To make $f_1(\theta) = f_2(\theta)$, we must set $A = 1$ and $B' = j$.

It follows that

$$e^{j\theta} = \cos \theta + j \sin \theta$$

This argument presumes the existence of a constant j whose square is -1 and that can be manipulated as an ordinary algebraic constant.

Where Does Euler's Formula Come From?

Euler's formula also follows from Maclaurin expansion of the exponential function, assuming the j behaves like any other algebraic constant.

Start with the expansion of the real-valued function:

$$e^{\theta} = 1 + \theta + \frac{\theta^2}{2!} + \frac{\theta^3}{3!} + \frac{\theta^4}{4!} + \frac{\theta^5}{5!} + \frac{\theta^6}{6!} + \frac{\theta^7}{7!} + \dots$$

Assume that the same expansion holds for complex-valued arguments:

$$\begin{aligned} e^{j\theta} &= 1 + j\theta + \frac{j^2\theta^2}{2!} + \frac{j^3\theta^3}{3!} + \frac{j^4\theta^4}{4!} + \frac{j^5\theta^5}{5!} + \frac{j^6\theta^6}{6!} + \frac{j^7\theta^7}{7!} + \dots \\ &= 1 + j\theta - \frac{\theta^2}{2!} - \frac{j\theta^3}{3!} + \frac{\theta^4}{4!} + \frac{j\theta^5}{5!} - \frac{\theta^6}{6!} - \frac{j\theta^7}{7!} + \dots \\ &= \underbrace{\left(1 - \frac{\theta^2}{2!} + \frac{\theta^4}{4!} - \frac{\theta^6}{6!} + \dots\right)}_{\cos \theta} + j \underbrace{\left(\theta - \frac{\theta^3}{3!} + \frac{\theta^5}{5!} - \frac{\theta^7}{7!} + \dots\right)}_{\sin \theta} \end{aligned}$$

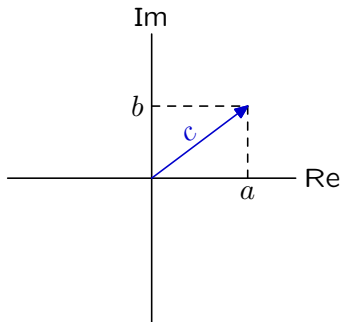
Euler's formula results by splitting the even and odd powers of θ .

$$e^{j\theta} = \cos \theta + j \sin \theta$$

Geometric Interpretation

In 1799, Caspar Wessel was the first to describe complex numbers as points in the complex plane. Imaginary numbers had been in use since the 1500's.

$$c = a + jb$$



Algebraic Addition

Addition: the real part of a sum is the sum of the real parts, and the imaginary part of a sum is the sum of the imaginary parts.

Let c_1 and c_2 represent complex numbers:

$$c_1 = a_1 + jb_1$$

$$c_2 = a_2 + jb_2$$

Then

$$c_1 + c_2 = (a_1 + jb_1) + (a_2 + jb_2) = (a_1 + a_2) + j(b_1 + b_2)$$

Geometric Addition

Rules for adding complex numbers are same as those for adding vectors.

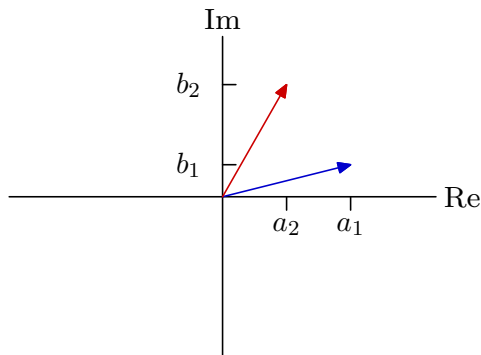
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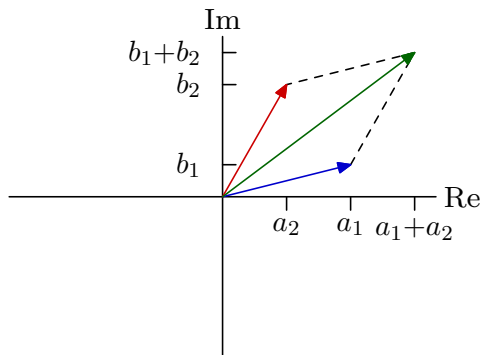
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Algebraic Multiplication

Multiplication is more complicated.

Let c_1 and c_2 represent complex numbers:

$$c_1 = a_1 + jb_1$$

$$c_2 = a_2 + jb_2$$

Then

$$\begin{aligned}c_1 \times c_2 &= (a + jb) \times (c + jd) \\ &= a \times c + a \times jd + jb \times c + jb \times jd \\ &= (ac - bd) + j(ad + bc)\end{aligned}$$

Although the rules of algebra apply, the result is complicated:

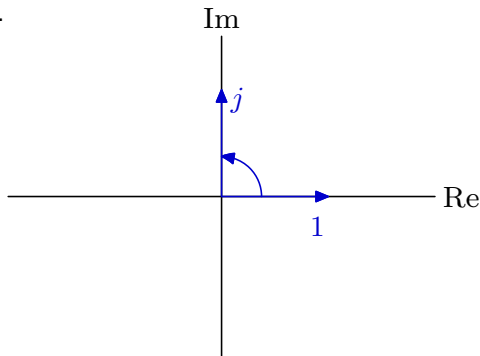
- the real part of a product is NOT the product of the real parts, and
- the imaginary part is NOT the product of the imaginary parts.

Geometric Multiplication

The two-dimensional view of complex numbers allows us to think about multiplication by an imaginary number as a **rotation**.

Multiplying by j

- **rotates 1 to j** ,
- rotates j to -1 ,
- rotates -1 to $-j$, and
- rotates $-j$ to 1.

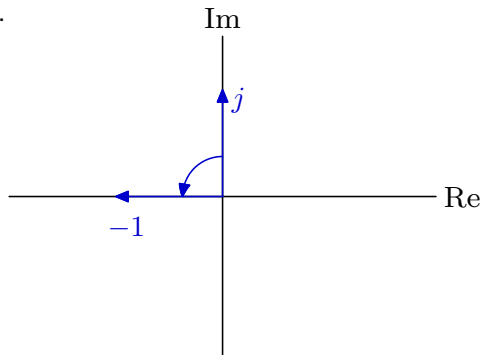


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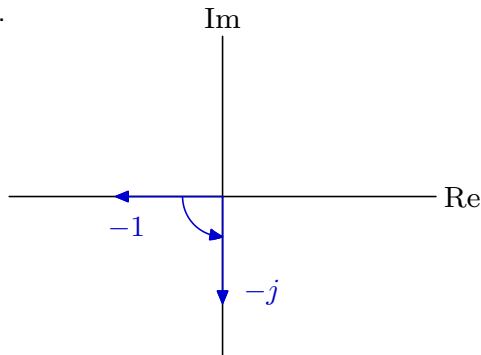


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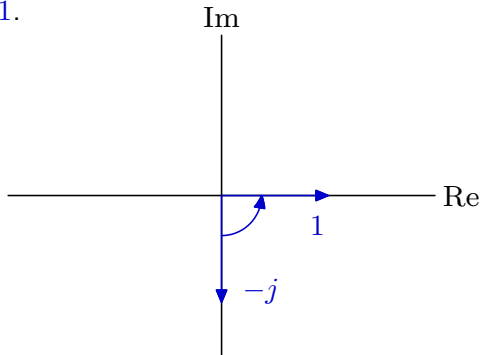


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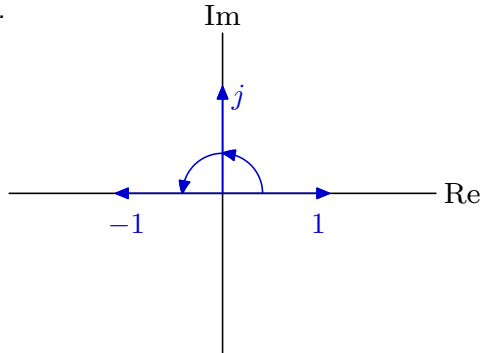


Geometric Multiplication

The two-dimensional view of complex numbers allows us to think about multiplication by an imaginary number as a **rotation**.

Multiplying by j

- rotates 1 to j ,
- rotates j to -1 ,
- rotates -1 to $-j$, and
- rotates $-j$ to 1.



Multiplying by j rotates a vector by $\pi/2$.

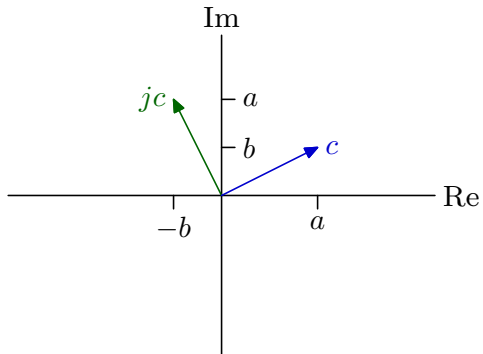
Multiplying by $j^2 = -1$ rotates a vector by π .

Geometric Multiplication

Multiplying by j rotates an arbitrary complex number by $\pi/2$.

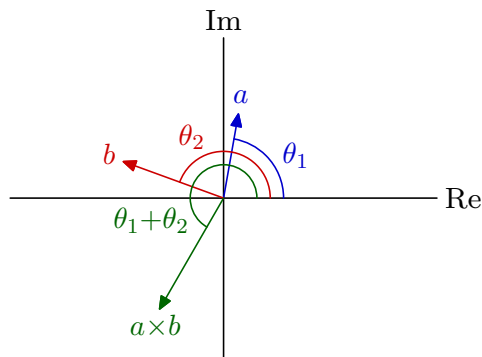
$$c = a + jb$$

$$jc = ja - b$$



Geometric Approach: Polar Form

The magnitude of the product of complex numbers is the **product** of their magnitudes. The angle of a product is the **sum** of the angles.

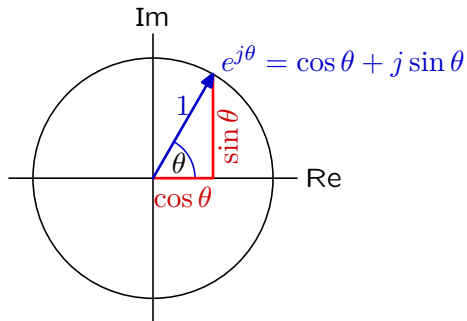


$$\begin{aligned}(r_1 \angle \theta_1) \times (r_2 \angle \theta_2) &= r_1(\cos \theta_1 + j \sin \theta_1) \times r_2(\cos \theta_2 + j \sin \theta_2) \\ &= r_1 r_2 \underbrace{(\cos \theta_1 \cos \theta_2 - \sin \theta_1 \sin \theta_2)}_{\cos(\theta_1 + \theta_2)} + j \underbrace{(\cos \theta_1 \sin \theta_2 + \sin \theta_1 \cos \theta_2)}_{\sin(\theta_1 + \theta_2)} \\ &= r_1 r_2 \angle (\theta_1 + \theta_2)\end{aligned}$$

Geometric Interpretation of Euler's Formula

Euler's formula equates polar and rectangular descriptions of a unit vector at angle θ .

$$e^{j\theta} = \cos \theta + j \sin \theta$$

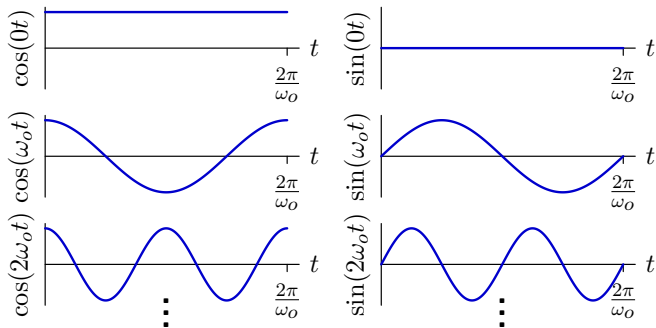


Simplifying Math By Using Complex Numbers

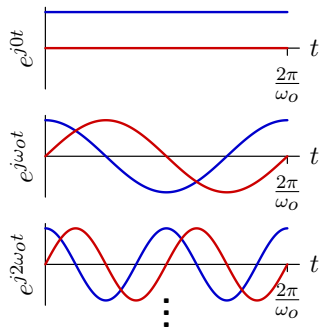
Euler's formula allows us to represent both sine and cosine basis functions with a single complex exponential:

$$f(t) = \sum \left(c_k \cos(k\omega_o t) + d_k \sin(k\omega_o t) \right) = \sum a_k e^{jk\omega_o t}$$

Real-valued basis functions



Complex basis functions



This halves the number of coefficients, but each is now complex-valued. More importantly, it replaces the trig functions with an exponential.

Converting From Trig Form To Complex Exponential Form

Assume that a function $f(t)$ can be written as a Fourier series in trig form.

$$f(t) = f(t + T) = c_0 + \sum_{k=1}^{\infty} \left(c_k \cos(k\omega_0 t) + d_k \sin(k\omega_0 t) \right)$$

We can use Euler's formula to convert sinusoids to complex exponentials.

$$e^{jk\omega_0 t} = \cos(k\omega_0 t) + j \sin(k\omega_0 t)$$

$$\cos(k\omega_0 t) = \operatorname{Re}\{e^{jk\omega_0 t}\} = (e^{jk\omega_0 t} + e^{-jk\omega_0 t})/2$$

$$\sin(k\omega_0 t) = \operatorname{Im}\{e^{jk\omega_0 t}\} = -j(e^{jk\omega_0 t} - e^{-jk\omega_0 t})/2$$

$$\begin{aligned} f(t) &= c_0 + \frac{1}{2} \sum_{k=1}^{\infty} \left(c_k e^{jk\omega_0 t} + c_k e^{-jk\omega_0 t} - j d_k e^{jk\omega_0 t} + j d_k e^{-jk\omega_0 t} \right) \\ &= c_0 + \frac{1}{2} \sum_{k=1}^{\infty} (c_k - j d_k) e^{jk\omega_0 t} + \frac{1}{2} \sum_{k=1}^{\infty} (c_k + j d_k) e^{-jk\omega_0 t} \\ &= c_0 + \frac{1}{2} \sum_{k=1}^{\infty} (c_k - j d_k) e^{jk\omega_0 t} + \frac{1}{2} \sum_{k=-1}^{-\infty} (c_{-k} + j d_{-k}) e^{jk\omega_0 t} = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} \end{aligned}$$

$$\text{where } a_k = \begin{cases} (c_k - j d_k)/2 & \text{if } k > 0 \\ c_0 & \text{if } k = 0 \\ (c_{-k} + j d_{-k})/2 & \text{if } k < 0 \end{cases}$$

Negative Frequencies

The complex form of a Fourier series has both positive and negative k 's.

Only positive values of k are used in the trig form:

$$f(t) = c_0 + \sum_{k=1}^{\infty} c_k \cos(k\omega_0 t) + \sum_{k=1}^{\infty} d_k \sin(k\omega_0 t)$$

but both positive and negative values of k are used in the exponential form:

$$f(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t}$$

If we only included positive k in the previous sum, the result would always have an imaginary component (unless $a_k = 0$ for all k).

If $f(t)$ is real-valued (as it must be for the trig form), then the complex coefficients a_k are **conjugate symmetric**:

$$a_{-k} = a_k^*$$

where the $*$ denotes the complex conjugate.

k :	...	-3	-2	-1	0	1	2	3	...
a_k :	...	$\frac{c_3 + jd_3}{2}$	$\frac{c_2 + jd_2}{2}$	$\frac{c_1 + jd_1}{2}$	c_0	$\frac{c_1 - jd_1}{2}$	$\frac{c_2 - jd_2}{2}$	$\frac{c_3 - jd_3}{2}$...

Fourier Series Directly From Complex Exponential Form

Assume that $f(t)$ is periodic in T and is composed of a weighted sum of harmonically related complex exponentials.

$$f(t) = f(t + T) = \sum_{k=-\infty}^{\infty} a_k e^{j\omega_0 k t}$$

We can “sift” out the component at $l\omega_0$ by multiplying both sides by $e^{-jl\omega_0 t}$ and integrating over a period.

$$\begin{aligned} \int_T f(t) e^{-j\omega_0 l t} dt &= \int_T \sum_{k=-\infty}^{\infty} a_k e^{j\omega_0 k t} e^{-j\omega_0 l t} dt = \sum_{k=-\infty}^{\infty} a_k \int_T e^{j\omega_0 (k-l)t} dt \\ &= \begin{cases} T a_l & \text{if } l = k \\ 0 & \text{otherwise} \end{cases} \end{aligned}$$

Solving for a_l provides an explicit formula for the coefficients:

$$a_k = \frac{1}{T} \int_T f(t) e^{-j\omega_0 k t} dt; \quad \text{where } \omega_0 = \frac{2\pi}{T}.$$

This formulation works even if $f(t)$ has complex values.

Orthogonality and Projection

Fourier components are separable because they are **orthogonal**.

Similar to separating a vector \bar{r} into x and y components.

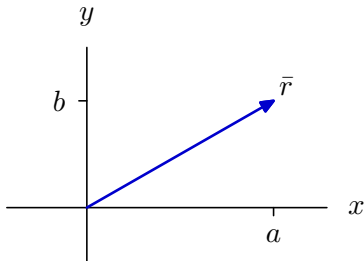
Since \hat{x} and \hat{y} are orthogonal, we can separate the x and y components of \bar{r} by **projection**:

$$a = \bar{r} \cdot \hat{x}$$

$$b = \bar{r} \cdot \hat{y}$$

Then

$$\bar{r} = a\hat{x} + b\hat{y}$$



Orthogonal Decompositions

Vector representation: let \bar{r} represent a vector with components a and b in the \hat{x} and \hat{y} directions, respectively.

$$a = \bar{r} \cdot \hat{x}$$

$$b = \bar{r} \cdot \hat{y}$$

(“analysis” equations)

$$\bar{r} = a\hat{x} + b\hat{y}$$

(“synthesis” equation)

Orthogonal Decompositions

Vector representation: let \bar{r} represent a vector with components a and b in the \hat{x} and \hat{y} directions, respectively.

$$a = \bar{r} \cdot \hat{x} \quad (\text{"analysis" equations})$$

$$b = \bar{r} \cdot \hat{y}$$

$$\bar{r} = a\hat{x} + b\hat{y} \quad (\text{"synthesis" equation})$$

Fourier series: let $f(t)$ represent a signal with harmonic components a_0, a_1, \dots, a_k for harmonics $e^{j0t}, e^{j\frac{2\pi}{T}t}, \dots, e^{j\frac{2\pi}{T}kt}$ respectively.

$$a_k = \frac{1}{T} \int_T f(t) e^{-j\frac{2\pi}{T}kt} dt \quad (\text{"analysis" equation})$$

$$f(t) = f(t + T) = \sum_{k=-\infty}^{\infty} a_k e^{j\frac{2\pi}{T}kt} \quad (\text{"synthesis" equation})$$

Orthogonal Decompositions

Integrating over a period **sifts** out the k^{th} component of the series.

Sifting as a dot product:

$$x = \bar{r} \cdot \hat{x} \equiv |\bar{r}| |\hat{x}| \cos \theta$$

Sifting as an inner product:

$$a_k = e^{j\frac{2\pi}{T}kt} \cdot f(t) \equiv \frac{1}{T} \int_T f(t) e^{-j\frac{2\pi}{T}kt} dt$$

where

$$a(t) \cdot b(t) = \frac{1}{T} \int_T a^*(t) b(t) dt.$$

The complex conjugate (*) makes the inner product of the k^{th} and m^{th} components equal to 1 iff $k = m$:

$$\frac{1}{T} \int_T \left(e^{j\frac{2\pi}{T}kt} \right)^* \left(e^{j\frac{2\pi}{T}mt} \right) dt = \frac{1}{T} \int_T e^{-j\frac{2\pi}{T}kt} e^{j\frac{2\pi}{T}mt} dt = \begin{cases} 1 & \text{if } k = m \\ 0 & \text{otherwise} \end{cases}$$

Check Yourself

How many of the following pairs of functions are orthogonal (\perp) in $T = 3$?

1. $\cos 2\pi t \perp \sin 2\pi t$?
2. $\cos 2\pi t \perp \cos 4\pi t$?
3. $\cos 2\pi t \perp \sin \pi t$?
4. $\cos 2\pi t \perp e^{j2\pi t}$?

Check Yourself

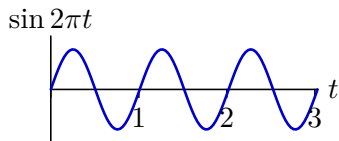
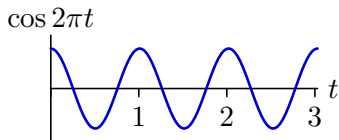
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$$\cos 2\pi t \perp \sin 2\pi t ?$$

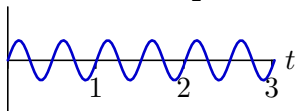
Check Yourself

How many of the following are orthogonal (\perp) in $T = 3$?

$\cos 2\pi t \perp \sin 2\pi t$?



$$\cos 2\pi t \sin 2\pi t = \frac{1}{2} \sin 4\pi t$$



$$\int_0^3 dt = 0 \text{ therefore YES}$$

Check Yourself

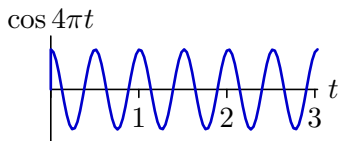
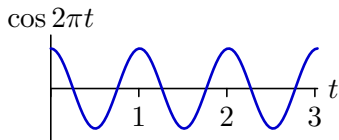
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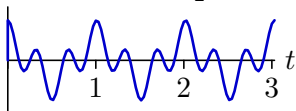
Check Yourself

How many of the following are orthogonal (\perp) in $T = 3$?

$\cos 2\pi t \perp \cos 4\pi t$?



$$\cos 2\pi t \cos 4\pi t = \frac{1}{2} \cos 6\pi t + \frac{1}{2} \cos 2\pi t$$



$$\int_0^3 dt = 0 \text{ therefore YES}$$

Check Yourself

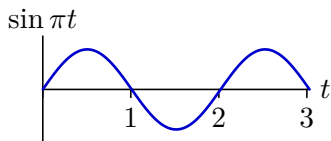
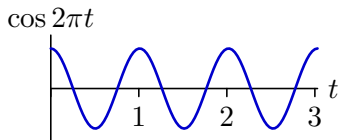
How many of the following are orthogonal (\perp) in $T = 3$?

$\cos 2\pi t \perp \sin \pi t$?

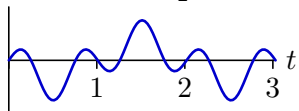
Check Yourself

How many of the following are orthogonal (\perp) in $T = 3$?

$\cos 2\pi t \perp \sin \pi t$?



$$\cos 2\pi t \sin \pi t = \frac{1}{2} \sin 3\pi t - \frac{1}{2} \sin \pi t$$



$$\int_0^3 dt \neq 0 \text{ therefore NO}$$

Check Yourself

How many of the following are orthogonal (\perp) in $T = 3$?

$$\cos 2\pi t \perp e^{j2\pi t} ?$$

Check Yourself

How many of the following are orthogonal (\perp) in $T = 3$?

$\cos 2\pi t \perp e^{j2\pi t}$?

$$e^{2\pi t} = \cos 2\pi t + j \sin 2\pi t$$

$\cos 2\pi t \perp \sin 2\pi t$ but not $\cos 2\pi t$

Therefore **NO**

Check Yourself

How many of the following pairs of functions are orthogonal (\perp) in $T = 3$? **2**

1. $\cos 2\pi t \perp \sin 2\pi t$? ✓

2. $\cos 2\pi t \perp \cos 4\pi t$? ✓

3. $\cos 2\pi t \perp \sin \pi t$? ✗

4. $\cos 2\pi t \perp e^{j2\pi t}$? ✗

Fourier Series

Comparison of trigonometric and complex exponential forms.

Complex Exponential Form

$$f(t) = f(t + T) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t}$$

$$a_k = \frac{1}{T} \int_T f(t) e^{-jk\omega_0 t} dt$$

Trigonometric Form

$$f(t) = f(t + T) = c_0 + \sum_{k=1}^{\infty} c_k \cos(k\omega_0 t) + \sum_{k=1}^{\infty} d_k \sin(k\omega_0 t)$$

$$c_0 = \frac{1}{T} \int_T f(t) dt$$

$$c_k = \frac{2}{T} \int_T \cos(k\omega_0 t) dt; \quad k = 1, 2, 3, \dots$$

$$d_k = \frac{2}{T} \int_T \sin(k\omega_0 t) dt; \quad k = 1, 2, 3, \dots$$

Comparison of Trigonometric and Complex Exponential Forms

It seems as though it takes more numbers to characterize the complex exponential form:

- Each harmonic frequency in the complex exponential form depends on two complex-valued numbers: a_k and a_{-k} .
- Each harmonic frequency in the trig form depends on two real-valued numbers: c_k and d_k .

Q: What is going on?

A: The complex exponential form allows $f(t)$ to have complex values. The trigonometric form requires that $f(t)$ be real-valued.

Q: Isn't it twice the work to compute both a_k and a_{-k} ?

A: Only if $f(t)$ is complex-valued.

If $f(t)$ is real-valued, then a_{-k} is the complex conjugate of a_k .

Is the Complex Exponential Form Actually Easier?

Last time, we determined the effect of a half-period shift on the Fourier coefficients of the trig form. The result was a bit complicated.

Assume that $f(t)$ is periodic in time with period T :

$$f(t) = f(t + T).$$

Let $g(t)$ represent a version of $f(t)$ shifted by half a period:

$$g(t) = f(t - T/2).$$

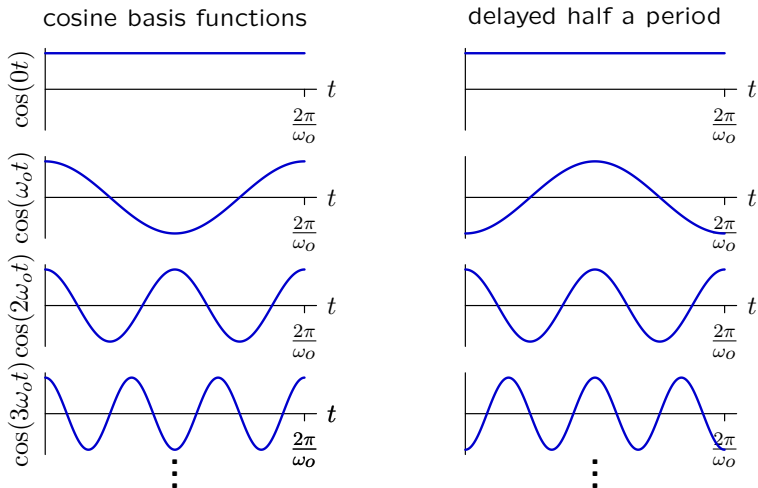
How many of the following statements correctly describe the effect of this shift on the Fourier series coefficients.

- cosine coefficients c_k are negated ✗
- sine coefficients d_k are negated ✗
- odd-numbered coefficients $c_1, d_1, c_3, d_3, \dots$ are negated ✓
- sine and cosine coefficients are swapped: $c_k \rightarrow d_k$ and $d_k \rightarrow c_k$ ✗

Half-Period Shift

Shifting $f(t)$ shifts the underlying basis functions of its Fourier expansion.

$$f(t - T/2) = \sum_{k=0}^{\infty} c_k \cos(k\omega_0(t - T/2)) + \sum_{k=1}^{\infty} d_k \sin(k\omega_0(t - T/2))$$

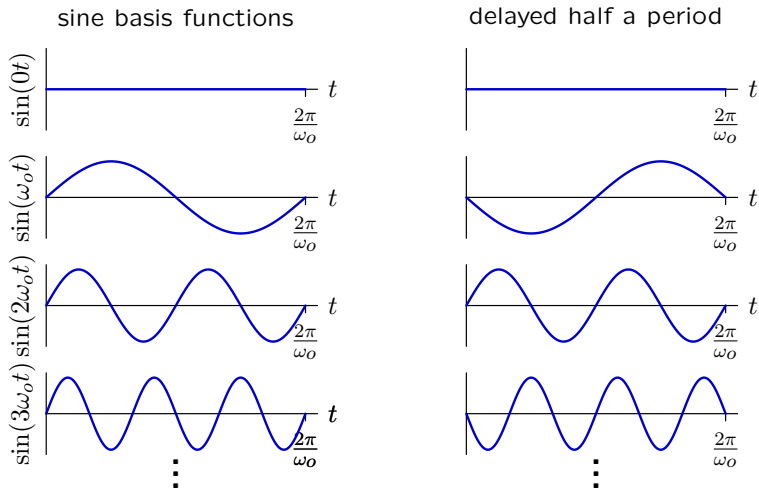


Half-period shift inverts c_k terms if k is odd. It has no effect if k is even.

Half-Period Shift

Shifting $f(t)$ shifts the underlying basis functions of its Fourier expansion.

$$f(t - T/2) = \sum_{k=0}^{\infty} c_k \cos(k\omega_0(t - T/2)) + \sum_{k=1}^{\infty} d_k \sin(k\omega_0(t - T/2))$$



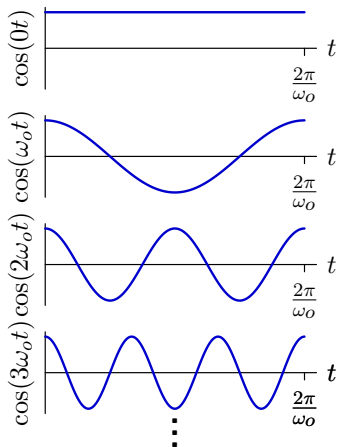
Half-period shift inverts d_k terms if and only if k is odd.

Quarter-Period Shift

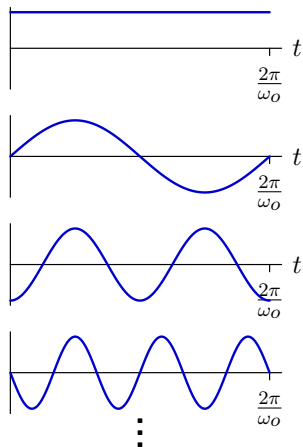
Shifting by $T/4$ is **even more complicated**.

$$f(t - T/4) = \sum_{k=0}^{\infty} c_k \cos(k\omega_o(t - T/4)) + \sum_{k=1}^{\infty} d_k \sin(k\omega_o(t - T/4))$$

cosine basis functions



delayed one fourth period



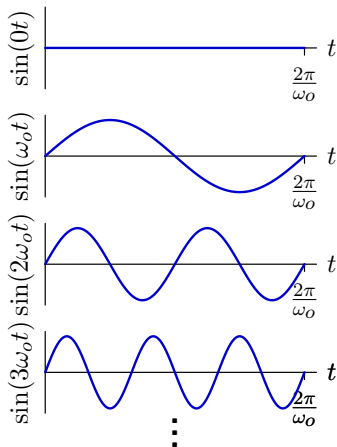
$$\cos(\omega_o t) \rightarrow \sin(\omega_o t); \quad \cos(2\omega_o t) \rightarrow -\cos(2\omega_o t); \quad \cos(3\omega_o t) \rightarrow -\sin(3\omega_o t)$$

Check Yourself: Alternative (more intuitive) Approach

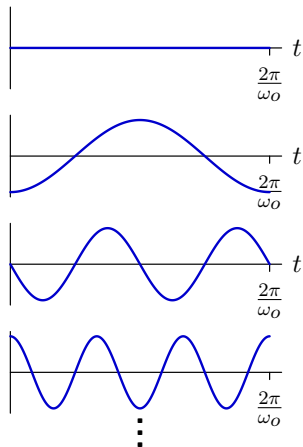
Shifting $f(t)$ shifts the underlying basis functions of its Fourier expansion.

$$f(t - T/4) = \sum_{k=0}^{\infty} c_k \cos(k\omega_o(t - T/4)) + \sum_{k=1}^{\infty} d_k \sin(k\omega_o(t - T/4))$$

sine basis functions



delayed 1/4 period



$$\sin(\omega_o t) \rightarrow -\cos(\omega_o t); \quad \sin(2\omega_o t) \rightarrow -\sin(2\omega_o t); \quad \sin(3\omega_o t) \rightarrow \cos(3\omega_o t)$$

Summary of Shift Results

Let c_k and d_k represent the Fourier series coefficients for $f(t)$

$$f(t) = f(t + T) = c_0 + \sum_{k=1}^{\infty} c_k \cos(k\omega_0 t) + \sum_{k=1}^{\infty} d_k \sin(k\omega_0 t)$$

and c'_k and d'_k represent those for a half-period delay.

$$g(t) = f(t - T/2) = c_0 + \sum_{k=1}^{\infty} c'_k \cos(k\omega_0 t) + \sum_{k=1}^{\infty} d'_k \sin(k\omega_0 t)$$

Then $c'_k = (-1)^k c_k$ and $d'_k = (-1)^k d_k$.

Let c''_k and d''_k represent those for a quarter-period delay.

$$g(t) = f(t - T/2) = c_0 + \sum_{k=1}^{\infty} c'_k \cos(k\omega_0 t) + \sum_{k=1}^{\infty} d'_k \sin(k\omega_0 t)$$

Then

$$c''_k = \begin{cases} c_k & \text{if } k = 0, 4, 8, 12, \dots \\ d_k & \text{if } k = 1, 5, 9, 13, \dots \\ -c_k & \text{if } k = 2, 6, 10, 14, \dots \\ -d_k & \text{if } k = 3, 7, 11, 15, \dots \end{cases} \quad d''_k = \begin{cases} d_k & \text{if } k = 0, 4, 8, 12, \dots \\ -c_k & \text{if } k = 1, 5, 9, 13, \dots \\ -d_k & \text{if } k = 2, 6, 10, 14, \dots \\ c_k & \text{if } k = 3, 7, 11, 15, \dots \end{cases}$$

Other Shifts Yield Even More Complicated Results

Let c_k and d_k represent the Fourier series coefficients for $f(t)$

$$f(t) = f(t + T) = c_0 + \sum_{k=1}^{\infty} c_k \cos(k\omega_0 t) + \sum_{k=1}^{\infty} d_k \sin(k\omega_0 t)$$

and c_k''' and d_k''' represent those for n eighth-period delay.

$$g(t) = f(t - T/8) = c_0 + \sum_{k=1}^{\infty} c_k' \cos(k\omega_0 t) + \sum_{k=1}^{\infty} d_k' \sin(k\omega_0 t)$$

$$c_k''' = \begin{cases} c_k & \text{if } k = 0, 8, 16, 24, \dots \\ \frac{\sqrt{2}}{2}(c_k + d_k) & \text{if } k = 1, 9, 17, 25, \dots \\ d_k & \text{if } k = 2, 10, 18, 26, \dots \\ \frac{\sqrt{2}}{2}(-c_k + d_k) & \text{if } k = 3, 11, 19, 27, \dots \\ -c_k & \text{if } k = 4, 12, 20, 28, \dots \\ \frac{\sqrt{2}}{2}(-c_k - d_k) & \text{if } k = 5, 13, 21, 29, \dots \\ -d_k & \text{if } k = 6, 14, 22, 30, \dots \\ \frac{\sqrt{2}}{2}(c_k - d_k) & \text{if } k = 7, 15, 23, 31, \dots \end{cases} \quad d_k''' = \dots$$

Effects of Time Shifts on Complex Exponential Series

Delaying time by τ multiplies the complex exponential coefficients of a Fourier series by a constant $e^{-jk\omega_0\tau}$.

Let a_k represent the complex exponential series coefficients of $f(t)$ and a'_k represent the complex exponential series coefficients of $g(t) = f(t - \tau)$.

$$\begin{aligned}a'_k &= \frac{1}{T} \int_T g(t) e^{-jk\omega_0 t} dt \\&= \frac{1}{T} \int_T f(t - \tau) e^{-jk\omega_0 t} dt \\&= \frac{1}{T} \int_T f(s) e^{-jk\omega_0 (s + \tau)} ds \\&= e^{-jk\omega_0 \tau} \frac{1}{T} \int_T f(s) e^{-jk\omega_0 s} ds \\&= e^{-jk\omega_0 \tau} a_k\end{aligned}$$

Each coefficient a'_k in the series for $g(t)$ is a constant $e^{-jk\omega_0 \tau}$ times the corresponding coefficient a_k in the series for $f(t)$.

Summary

We introduced the complex exponential form of Fourier series.

- complex numbers
- complex exponentials and their relation to sinusoids
- analysis and synthesis with complex exponentials
- delay property: much simpler with complex exponentials

Trig Table

$$\sin(a+b) = \sin(a) \cos(b) + \cos(a) \sin(b)$$

$$\sin(a-b) = \sin(a) \cos(b) - \cos(a) \sin(b)$$

$$\cos(a+b) = \cos(a) \cos(b) - \sin(a) \sin(b)$$

$$\cos(a-b) = \cos(a) \cos(b) + \sin(a) \sin(b)$$

$$\tan(a+b) = (\tan(a)+\tan(b))/(1-\tan(a) \tan(b))$$

$$\tan(a-b) = (\tan(a)-\tan(b))/(1+\tan(a) \tan(b))$$

$$\sin(A) + \sin(B) = 2 \sin((A+B)/2) \cos((A-B)/2)$$

$$\sin(A) - \sin(B) = 2 \cos((A+B)/2) \sin((A-B)/2)$$

$$\cos(A) + \cos(B) = 2 \cos((A+B)/2) \cos((A-B)/2)$$

$$\cos(A) - \cos(B) = -2 \sin((A+B)/2) \sin((A-B)/2)$$

$$\sin(a+b) + \sin(a-b) = 2 \sin(a) \cos(b)$$

$$\sin(a+b) - \sin(a-b) = 2 \cos(a) \sin(b)$$

$$\cos(a+b) + \cos(a-b) = 2 \cos(a) \cos(b)$$

$$\cos(a+b) - \cos(a-b) = -2 \sin(a) \sin(b)$$

$$2 \cos(A) \cos(B) = \cos(A-B) + \cos(A+B)$$

$$2 \sin(A) \sin(B) = \cos(A-B) - \cos(A+B)$$

$$2 \sin(A) \cos(B) = \sin(A+B) + \sin(A-B)$$

$$2 \cos(A) \sin(B) = \sin(A+B) - \sin(A-B)$$